

6.1
(p. 149)

Topographic Rossby Waves

linear, hydrostatic, Boussinesq Approx.:

(1) $u_t - fv = -\frac{1}{\rho_0} p_x$

(2) $v_t + fu = -\frac{1}{\rho_0} p_y$

(3) $0 = -\frac{1}{\rho} p_z - \frac{g\rho}{\rho_0}$

→ (4) $p_t + w\rho_0 z = 0$ } → $w = -\frac{1}{\rho_0 N^2} p_{zt}$

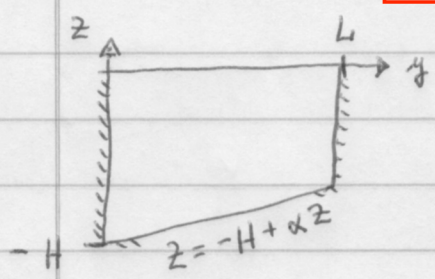
(5) $u_x + v_y + w_z = \sigma$

WAVE EQUATION

$$\left[p_{xx} + p_{yy} + \left(\frac{\partial^2}{\partial t^2} + f^2 \right) \left(\frac{p_z}{N^2} \right)_z \right] = \sigma$$

$e^{-i\sigma t}$:

$$p_{xx} + p_{yy} + (f^2 - \sigma^2) \left(\frac{p_z}{N^2} \right)_z = \sigma$$



- $v = 0$ @ $y = 0$ and $y = L$
- $w = 0$ @ $z = 0$
- $w = \alpha v$ @ $z = -H + \alpha z$

bottom slope α

$$y = 0, L \quad v = 0 \quad \downarrow \quad \underline{i\sigma p_y + f p_x = 0}$$

$$z = 0 \quad w = 0 \quad \downarrow \quad \underline{p_z = 0}$$

$$z = -H + \alpha z \quad w = \alpha v \quad \downarrow \quad \underline{i\sigma(f^2 - \sigma^2)p_z = \alpha N^2(i\sigma p_y + f p_x)}$$

W V

Scale (x, y) by L

z by H

Set $w = \sigma/f$

$N = \text{const.}$

Then our wave problem with BC becomes

$$p_{xx} + p_{yy} + \underbrace{\frac{(1-w^2)}{S^2}}_{S^2} p_{zz} = 0$$

$$i w p_y + p_x = 0 \quad \text{at } y = 0, 1$$

$$p_z = 0 \quad \text{at } z = 0$$

Bottom Boundary Condition:

$$i w (1-w^2) p_z = \delta S^2 (i w p_y + p_x) \quad \text{at } z = 1 + \delta y$$

where $\delta = \frac{\alpha L}{H}$ is scaled bottom slope

Internal
Rossby radius

$$S' = (N \cdot H) / (f L)$$

is Burger number $S' = \frac{N \cdot H}{f} / L$

Burger number S' measures the importance of stratification relative to the spatial scale of the

motion
$$S' = \frac{\text{internal Rossby radius}}{\text{geometric length scale}} = \frac{N \cdot H / f}{L}$$

Large S' means strong stratification $L_D > L$
(and/or large aspect ratio H/L)

Small S' means weak stratification $L_D < L$
(and/or nearly horizontal motions)

Consider $\omega \ll 1$ ($\sigma \ll f$) low-frequency
 $\delta \ll 1$ (gentle bottom slope)

$$p_{xx} + p_{yy} + \frac{1}{S'^2} p_{zz} = \sigma$$

$$y = 0, 1 : p_x = \sigma$$

$$z = \sigma : p_z = \sigma$$

$$z = -1 : i \omega p_z = \delta S'^2 p_x$$

Solution is

$$p = e^{ikx} \sin(n\pi y) \cosh(\mu z)$$

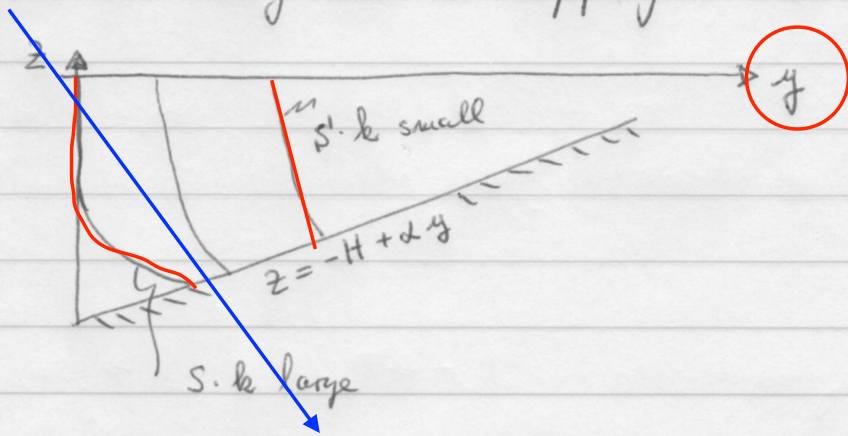
where

$$\mu^2 = S'^2 (n^2 + k^2)$$

gives a vertical decay scale

Thus strong stratification and/or short spatial scales

result in strong bottom trapping



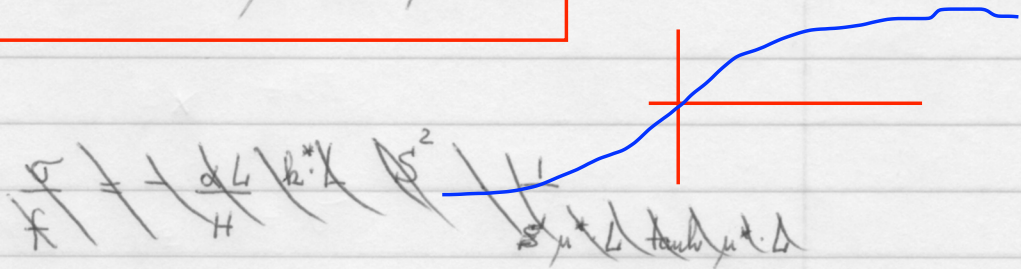
Dispersion relation becomes (non-dimensional)

$$\omega = - \frac{\delta k S'^2}{\mu \tanh \mu}$$

k is the across-slope component of the wave number vector

Compare this to the planetary Rossby Wave Dispersion

propagates with shallow water on the right (northern Hemisphere)



For weak stratification $\sigma' \rightarrow \sigma$

$\downarrow \mu \rightarrow \sigma$

$\delta \mu \approx \sigma' k \mu \rightarrow \sigma'^2 = \sigma'^2 (n^2 \pi^2 + k^2)$

and

$$\omega = - \frac{\delta k \sigma'^2}{\sigma'^2 (n^2 \pi^2 + k^2)} = - \frac{\delta k}{n^2 \pi^2 + k^2}$$

barotropic topographic Rossby Waves

or dimensionally

$$\sigma = - \frac{\alpha}{H} k \cdot f \cdot \frac{1}{[(n\pi/L)^2 + k^2]}$$

k is the along-slope wave number vector component

This is the dispersion for topographic Rossby Waves

$$\sigma = - \frac{\alpha k f}{H [(n\pi/L)^2 + k^2]}$$

$$\sigma = - \frac{f \alpha k}{H (l^2 + k^2)}$$

$$K^2 = l^2 + k^2 \gg \frac{f^2}{gH}$$

horizontal wave number length much smaller than Rossby radius

$$\frac{2\pi}{K} \ll \frac{\sqrt{gH}}{f}$$

short waves